

PHY208 – atoms and lasers Lecture 4

Introduction to spectroscopy

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What is spectroscopy ?

Analysis of the energy distribution of a (light) signal

What fraction of the signal is carried by modes with energy between hv and h(v+dv) ?

What is the ability of the system under scrutiny to emit (or absorb) light at a given wavelength ?







The atomic barcode



Spectral analysis of sunlight

Fraunhofer lines 1814

Norman Lockyer 1868



Introduction

The atomic barcode



TIM Cu Sr Ba Na Ca Li Barlum Copper Calcium Lithium Sodium Strontium

Practical applications

sciencenotes.org

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Probing the Universe – Lyman-alpha forest



Spectrum from nearby star

Spectrum from a (very) distant quasar





Outline of lecture 4

I. Light emitting transitions

Which light is actually emitted by a sample ?

II. Line width : broadening effects

What is the detailed shape of an atomic line? How is this shape affected by the environment ?

III. Line shift : the Zeeman effect

How does the environnement shifts the emissions ? How does the emissions probe the environnement ?

Not addressed in this lecture : experimental techniques





Optical transitions

What light is actually emitted ?

What are the atomic states ? (Transition from one state to another)

How are these states populated ?

(Absorption from an occupied low energy state to an empty high energy state.) Emission from an occupied high energy state to an empty low energy state.)

What light correspond to this transition ? (No mode, no transition)









Energy states

Internal structure

Attractive Coulomb potential from nucleus

Quantitative treatment in lecture 6

Repulsive interaction from the other electrons

(Interaction due to the spin of the electron)







Example: atomic states for Hydrogen

Results from the most simple Schrodinger model (full derivation in lecture 6)

n=3 ______









I. Light emitting transition

1 = 0

n=1

Example: atomic states for Hydrogen

Results from the most simple Schrodinger model (full derivation in lecture 6)



Energy levels :

Total angular momentum :

Label : S (l=0), P (l=1), D(l=2), f...



 $E_n = -\frac{1}{n^2} \frac{m_e e^4}{2\hbar^2}$



Example: atomic states for Hydrogen

Results from the most simple Schrodinger model (full derivation in lecture 6)

 $E_n = -\frac{1}{n^2} \frac{m_e e^4}{2\hbar^2}$ Energy levels : -1 0 1 -2 -1 0 1 2 m = 0 $L^2 = l(l+1)\hbar^2$ Total angular momentum : -1 0 1 m = 0Label : S (l=0), P (l=1), D(l=2), f...Angular momentum along z : $L_z = m\hbar$

Each energy state is degenerated n² times.



 $l \leq n-1$





n=3

n=2

Energy states

Internal structure

I. Light emitting transition

Attractive Coulomb potential from nucleus Quantitative treatment in lecture 6 $H_{\rm atom}$ **Energy levels** Repulsive interaction from the other electrons for the atom "alone" (Interaction due to the spin of the electron) (n, L, m, S...) **External environment** $H_{\rm env.}$ Magnetic field Influence of **Electric field** external parameters Light (**R**, **P**) Interaction with other atoms





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Light: allowed and forbidden transitions

Selection rule : basic conservation

(semi quantum or quantum model)

Energy $\hbar\omega = E_f - E_i$

Angular momentum

 $m_f - m_i = 0$ $m_f - m_i = \pm 1$

Before Atom in an excited state Atom in a ground state 0 photon E_i

 $m_i\hbar$

 $m_f \hbar + L_{z, \text{photon}}$

After

1 photon

 $E_f + \hbar\omega$

 π polarization \rightarrow L₇ = 0 σ_+ polarization $\rightarrow L_7 = \pm \hbar$

Available modes

How many modes satisfy these selection rules ?



Energy

input





Populations: (whatever)-escence

Energy in, photon out.

Incandescence

Energy from thermal agitation

 $\frac{n(E+dE)}{n(E)} = e^{-\frac{dE}{kT}}$

Photo-luminescence

Energy from light absorption

Fluorescence

Phosphorescence

Electro-luminescence

Energy from electronic injection

Cathodo-luminescence

Energy from electron beam

I. Light emitting transition



input

Light

output



Kirchhoff law of radiation :

 $\alpha(\nu) = \epsilon(\nu) \quad \forall \nu$

(Note for the future)

How to calculate quantitatively these rates : Fermi Golden Rule

$$R(h\nu) \propto \sum_{i,j\neq i} |\langle \psi_j| - e\hat{\mathbf{r}}.\mathbf{E} |\psi_i\rangle|^2 \,\delta \left(E_j - E_i - h\nu\right) \, p_i \left(1 - p_j\right)$$

Full quantum expression for the spontaneous decay rate :

$$\Gamma_{\rm vac} = \frac{|\langle g|e\hat{\mathbf{r}}|e\rangle|^2 \omega_{eg}^3}{3\pi\epsilon_0 \hbar c^3}$$

I. Light emitting transition



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Intrinsic broadening





Comment on spontaneous emission

Consider an atom in an energy eigenstate of $H_{
m atom}$

$$i\hbar \frac{\partial}{\partial t}\psi = H_{\rm atom}\psi \Rightarrow \psi(t) = \psi_0 e^{-iEt/\hbar}$$

the atom remains in the eigenstate forever...

What actually causes spontaneous emission ? What gives the decay rate Γ ?

Collision between atoms

Collisionnal broadening

Nature of interactions, density (pressure)...

But what if the atom is alone ?

Need quantum description for the field.

$$\Gamma_{\rm vac} = \frac{\left|\langle g | \hat{\mathbf{D}} | e \rangle \right|^2 \omega_{eg}^3}{3\pi\epsilon_0 \hbar c^3}$$

 $\Gamma_{\rm coll}$

 $\simeq 1 - 100 \,\mathrm{MHz}$

for an allowed atomic transition

$$\Gamma = \Gamma_{\rm coll} + \Gamma_{\rm vac} + \dots$$



Thermal broadening



Previous model vs actual result for a bunch of atom...

- > Not a Lorentzian shape
- Much broader than Γ
- > Depends on the temperature...

Should include the thermal motion of the atoms !





The Doppler effect







Doppler Frequency shift :

$$\omega = \frac{\omega_0}{1 + \frac{V}{c}}$$

III. Linewidth

Doppler broadening

 $I_{\text{tot}}(\omega) = \int dv_x dv_y dv_z I_{\text{at}} \left(\omega \left(1 + \frac{v_x}{c} \right) \right) \, n_{\text{at}} \left(\mathbf{v} \right)$

 $n_{\rm at}(\mathbf{v}) \propto \exp\left(-\frac{\frac{1}{2}m(v_x^2+v_y^2+v_z^2)}{kT}\right)$

 $I_{\rm at}(\omega) \propto \delta(\omega - \omega_0)$

Thermal distribution

Considering sharp emission profile (atomic emission $@\omega_0$ only)

$$I_{\text{tot}}(\omega) \propto \int dv_x \,\delta\left(v_x - c\frac{\Delta\omega}{\omega}\right) \,n_{\text{at}}\left(\mathbf{v}\right)$$

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Thermal broadening:

$$I_{\text{tot}} \simeq \frac{I_0}{\sqrt{2\pi\Delta\omega^2}} \exp\left(-\frac{(\omega-\omega_0)^2}{2\Delta\omega^2}\right)$$









Voigt profile

Considering actual Lorentzian emission profile

Voigt profile (no analytical expression)

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$$\begin{split} I_{\rm tot}(\omega) \propto \exp\left(-\frac{1}{2}\frac{mc^2}{k_BT\omega_0^2}\left(\omega-\omega_0\right)^2\right)\\ \text{Gaussian Doppler profile}: \quad \Delta\omega = \sqrt{\frac{k_BT\omega_0^2}{mc^2}} \end{split}$$

Considering sharp emission profile (atomic emission
$$@\omega_0$$
 only)

$$I_{\text{tot}}(\omega) = \int dv_x dv_y dv_z I_{\text{at}} \left(\omega \left(1 + \frac{v_x}{c}\right)\right) n_{\text{at}} \left(\omega \left(1 + \frac{v_x}{c}\right)\right) = 0$$

$$t_{\text{tot}}(\omega) = \int dv_x dv_y dv_z I_{\text{at}} \left(\omega \left(1 + \frac{v_x}{c} \right) \right) \, n_{\text{at}} \left(\mathbf{v} \right)$$

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rp emission profile (atomic emission @
$$\omega_0$$
 only)
 $\sigma_{
m t}(\omega) \propto \exp\left(-rac{1}{2}rac{mc^2}{k_BT\omega_0^2}\left(\omega-\omega_0\right)^2
ight)$

Г

1.0



Voigt profile



$$I_{\text{tot}}(\omega) = \int dv_x dv_y dv_z I_{\text{at}}\left(\omega\left(1 + \frac{v_x}{c}\right)\right) n_{\text{at}}\left(\mathbf{v}\right)$$

Considering sharp emission profile (atomic emission $@\omega_0$ only)

$$I_{
m tot}(\omega) \propto \exp\left(-rac{1}{2}rac{mc^2}{k_B T \omega_0^2} \left(\omega - \omega_0
ight)^2
ight)$$

Gaussian Doppler profile : $\Delta \omega = \sqrt{rac{k_B T \omega_0^2}{mc^2}}$

Considering actual Lorentzian emission profile

Voigt profile (no analytical expression)

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Zeeman effect (classical)



 $m_e \frac{d^2}{dt^2} \mathbf{r} = -m_e \omega_0^2 \mathbf{r} - q \mathbf{v} \times (B_{\text{ext}} \mathbf{u}_{\mathbf{z}})$

Motion along z

 $\ddot{z} + m\omega_0^2 z = 0$

Resonant frequency :



Motion in the transverse plane

 $\begin{cases} \ddot{x} + \omega_0^2 x + \frac{qB}{m} y = 0\\ \ddot{y} + \omega_0^2 y - \frac{qB}{m} x = 0 \end{cases}$

Resonant frequency :



 $\mu_B = \frac{e\hbar}{2m_e} = 14 \,\mathrm{GHz/T}$

Corresponding to σ_{\pm} polarization



Zeeman effect (classical - continued)







Around 640nm (Cadmium emission line)

II. Lineshift : the Zeeman effect

Anomalous Zeeman effect

Let's try on another transition !

Around 546 nm (Mercury emission line)



II. Lineshift : the Zeeman effect

Credit : Dave Gardner University of Scranton





From classical to semi-classical



Atom without mag-field :

 $H_{\rm atom}$

How to include the influence of the mag field ?

(reminder from PHY 205)

Elecron orbiting around the proton \rightarrow electrical current

Electrical current \rightarrow magnetic dipole

Magnetic dipole \rightarrow magnetic energy

Magnetic energy \rightarrow Zeeman Hamiltonian

II. Lineshift : the Zeeman effect

$$I = \frac{-q}{T} = \frac{-q v}{2\pi r}$$
$$\mathcal{M} = I \mathbf{S} = \frac{-q r v}{2} \mathbf{u}_{\mathbf{z}} = \frac{-q}{2m_e} \mathbf{L}$$
$$E_p = -\mathcal{M} \cdot \mathbf{B} = \frac{q}{2m_e} \mathbf{L} \cdot \mathbf{B}$$
$$\hat{V} = \underbrace{\left(q \cdot \frac{q}{2m_e} \hat{\mathbf{L}}\right)}_{e_{2m_e}} \cdot \hat{\mathbf{B}} \cdot \mathbf{B}$$

g_e: Landé factor (derivation from full quantum atom)

μ

II. Lineshift : the Zeeman effect

Toolbox : perturbation theory

How does a small change in the Hamiltonian changes the energy levels ?

Without perturbation :

 $\hat{H}_0|\psi_n\rangle = E_n|\psi_n\rangle$

 $H_{\rm atom} + V$

$$\left(\hat{H}_0 + \hat{V}\right) |\psi\rangle = E |\psi\rangle$$

1st order development $|\psi
angle = |\psi_n
angle + |\delta\psi
angle + ...$

 $E = E_n + \Delta E + \dots$

Energy shift :
$$\Delta E_n = \langle \psi_n | \hat{V} | \psi_n \rangle$$

For the Zeeman effect :

$$\Delta E_n = \left(\frac{qB}{2m_e}g_e\right) \langle \psi_n | \hat{L_z} | \psi_n \rangle$$



n=3

n=2

n=1

Example: atomic states for Hydrogen

Results from the most simple Schrodinger model (full derivation in lecture 6)





Atomic states

Results from the most simple Schrodinger model (full derivation in lecture 6)



Without magnetic field, each energy state is degenerated n² times.

Mag. Field lifts the degeneracy : each energy shifts according to its angular momentum along z

$$\Delta E = \frac{qB}{2m_e} g_e \langle n, l, m | \hat{L_z} | n, l, m \rangle = \frac{qB}{2m_e} g_e m\hbar$$

II. Lineshift : the Zeeman effect

Cd-Hg Zeeman shifts













II. Lineshift : the Zeeman effect

Ex: magnetic sunspot





Ex: magnetic sunspot





The usual effects

Magnetic field

Zeeman shift

Stark shift

Lamp (or light) shift

